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discusses rotations
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dynamics. (Note that 4.3 is a linear algebra review. Other sections from chapter 4 and chapter 5 may be of interest, but the above are the most important ones.) 4:

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the body frame
unit vectors as $i, j,$**

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Hamiltonian is now
explicitly
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fact that $dH=dt6=$
0. The energy is
given by $E= T+ V =$
 $1/2 (Q_{\dot{}} +$
 $b^2 \cos^2(\omega t))^2 + 1/2 (k_1 + k_2)(Q +$
 $b \sin(\omega t))^2$ (55) So,
 $dE/dt = m(Q_{\dot{}} +$
 $b \cos(\omega t))(Q$
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5(Goldstein 2.19)

If the mass distribution has a given symmetry, so will the potential and therefore so will the Lagrangian. From the symmetry, we deduce the conserved quantity.

- (a) The force does not depend on**

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$(x,y) \rightarrow (p_x,p_y)$
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unit vectors as \mathbf{i} , \mathbf{j} ,
and \mathbf{k} , and**

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Hamiltonian is now
explicitly

dependent on time,
and hence is not
conserved, as is

confirmed by the
fact that $dH/dt \neq$

0. The energy is
given by $E = T + V =$

$\frac{1}{2} m \dot{Q}^2 +$

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$b^2 \cos^2(\omega t) + \frac{1}{2} (k_1 + k_2) (Q_0 + b \sin(\omega t))^2$ (55) So,
 $\frac{dE}{dt} = m \dot{Q}_0 + 2b \cos(\omega t) (Q_0 + b \sin(\omega t)) \omega + (k_1 + k_2) b \sin(\omega t) \omega$

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distribution has a

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given symmetry, so
will the potential
and therefore so
will the
Lagrangian. From
the symmetry, we
deduce the
conserved quantity.

- (a) The force
does not depend on
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